

Electromagnetic Plane Wave (1)



Maxwell's equations for electric and magnetic fields in free space (no sources):

- Gauss' laws: $\oint \vec{E} \cdot d\vec{A} = 0, \quad \oint \vec{B} \cdot d\vec{A} = 0.$
- Faraday's and Ampère's laws: $\oint \vec{E} \cdot d\vec{\ell} = -\frac{d\Phi_B}{dt}, \quad \oint \vec{B} \cdot d\vec{\ell} = \mu_0\epsilon_0 \frac{d\Phi_E}{dt}.$

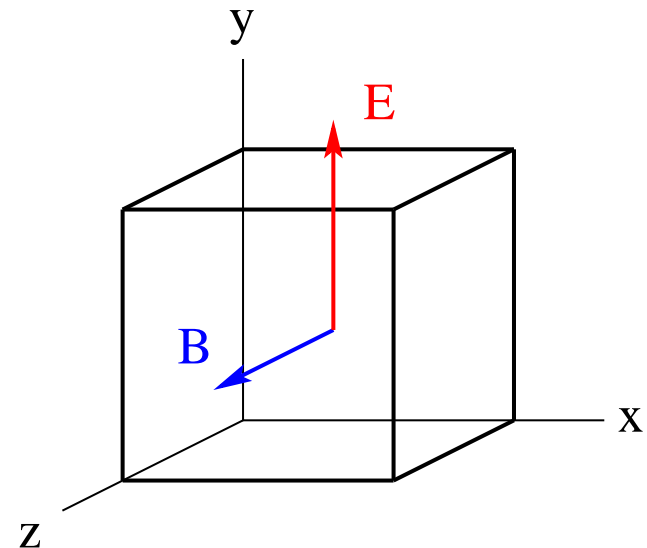
Consider fields of particular directions and dependence on space:

$$\vec{E} = E_y(x, t)\hat{j}, \quad \vec{B} = B_z(x, t)\hat{k}.$$

Gauss' laws are then automatically satisfied.

Use the cubic Gaussian surface to show that

- the net electric flux Φ_E is zero,
- the net magnetic flux Φ_B is zero.

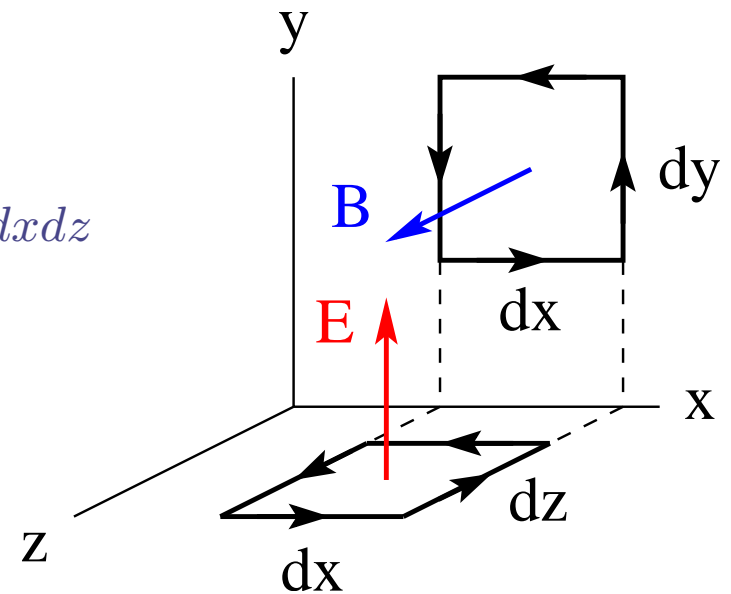


Electromagnetic Plane Wave (2)



- Faraday's law, $\oint \vec{E} \cdot d\vec{\ell} = -\frac{d\Phi_B}{dt}$,
applied to loop in (x, y) -plane becomes
 $[E_y(x + dx, t) - E_y(x, t)]dy = -\frac{\partial}{\partial t}B_z(x, t)dx dy$
 $\Rightarrow \frac{\partial}{\partial x}E_y(x, t) = -\frac{\partial}{\partial t}B_z(x, t) \quad (\text{F})$

- Ampère's law, $\oint \vec{B} \cdot d\vec{\ell} = \mu_0\epsilon_0 \frac{d\Phi_E}{dt}$,
applied to loop in (x, z) -plane becomes
 $[-B_z(x + dx, t) + B_z(x, t)]dz = \mu_0\epsilon_0 \frac{\partial}{\partial t}E_y(x, t)dx dz$
 $\Rightarrow -\frac{\partial}{\partial x}B_z(x, t) = \mu_0\epsilon_0 \frac{\partial}{\partial t}E_y(x, t) \quad (\text{A})$



Electromagnetic Plane Wave (3)



Take partial derivatives $\frac{\partial}{\partial x}$ (F) and $\frac{\partial}{\partial t}$ (A): $\frac{\partial^2 E_y}{\partial x^2} = -\frac{\partial^2 B_z}{\partial t \partial x}$, $-\frac{\partial^2 B_z}{\partial t \partial x} = \mu_0 \epsilon_0 \frac{\partial^2 E_y}{\partial t^2}$.

$$\Rightarrow \frac{\partial^2 E_y}{\partial t^2} = c^2 \frac{\partial^2 E_y}{\partial x^2} \quad (\text{E}) \quad (\text{wave equation for electric field}).$$

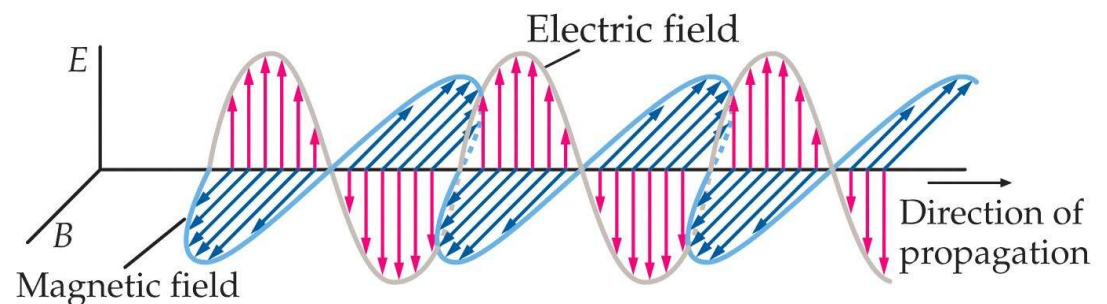
Take partial derivatives $\frac{\partial}{\partial t}$ (F) and $\frac{\partial}{\partial x}$ (A): $\frac{\partial^2 E_y}{\partial t \partial x} = -\frac{\partial^2 B_z}{\partial t^2}$, $-\frac{\partial^2 B_z}{\partial t^2} = \mu_0 \epsilon_0 \frac{\partial^2 E_y}{\partial t \partial x}$.

$$\Rightarrow \frac{\partial^2 B_z}{\partial t^2} = c^2 \frac{\partial^2 B_z}{\partial x^2} \quad (\text{B}) \quad (\text{wave equation for magnetic field}).$$

$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}} \quad (\text{speed of light}).$$

Sinusoidal solution:

- $E_y(x, t) = E_{max} \sin(kx - \omega t)$
- $B_z(x, t) = B_{max} \sin(kx - \omega t)$



Electromagnetic Plane Wave (4)



For given wave number k the angular frequency ω is determined, for example by substitution of $E_{max} \sin(kx - \omega t)$ into (E).

For given amplitude E_{max} the amplitude B_{max} is determined, for example, by substituting $E_{max} \sin(kx - \omega t)$ and $B_{max} \sin(kx - \omega t)$ into (A) or (F).

$$\Rightarrow \frac{\omega}{k} = \frac{E_{max}}{B_{max}} = c.$$

The direction of wave propagation is determined by the Poynting vector:

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}.$$

