

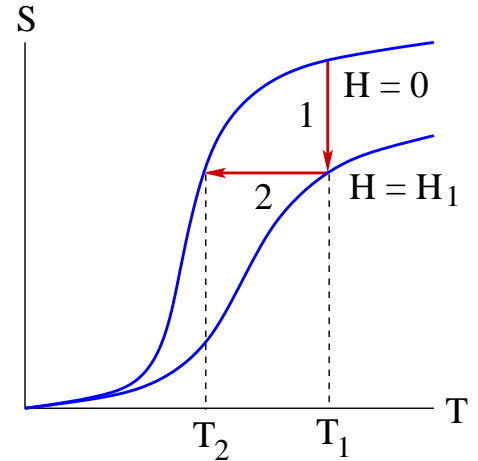
# Adiabatic demagnetization [tln24]

Equation of state for paramagnetic salt in a weak magnetic field:  $M(T, H) = \chi_T(T)H$ .

Helmholtz free energy:

$$dA = -SdT + HdM, \quad \left(\frac{\partial A}{\partial M}\right)_T = H = \frac{M}{\chi_T}$$

$$\Rightarrow A(T, M) = A(T, 0) + \frac{M^2}{2\chi_T}.$$



Entropy:  $S(T, M) = -\left(\frac{\partial A}{\partial T}\right)_M = S(T, 0) - \frac{1}{2}M^2 \left[\frac{d}{dT}\chi_T^{-1}\right]$

$$\Rightarrow S(T, H) = S(T, 0) + \frac{1}{2}H^2 \frac{d\chi_T}{dT}; \quad \chi_T > 0, \quad \frac{d\chi_T}{dT} < 0 \text{ for paramagnet.}$$

Third law:  $\lim_{T \rightarrow 0} S(T, H) = 0$  independent of  $H \Rightarrow \lim_{T \rightarrow 0} \frac{d\chi_T}{dT} = 0$ .

1. Isothermal magnetization:  $\Delta S = \frac{1}{2} \frac{d\chi_T}{dT} H_1^2 < 0$ .

Heat expelled from system:  $\Delta Q = T_1 \Delta S$ .

2. Adiabatic demagnetization:  $\Delta S = 0 \Rightarrow S(T_1, H_1) = S(T_2, 0)$ .

$$\Rightarrow S(T_2, 0) = S(T_1, 0) + \frac{1}{2}H_1^2 \frac{d\chi_T}{dT} \Big|_{T_1} \Rightarrow T_2 < T_1.$$

Consider the entropy function  $S(T, H)$  for iron ammonium alum.

- The sequence of steps 1 and 2 approaches absolute zero. As  $T \rightarrow 0$ , adiabates and isotherms become increasingly parallel, implying a diminishing efficiency of the cooling process.
- The sequence of steps requires heat reservoirs at various temperatures. They can be established by employing a Carnot engine consisting of steps 1 and 2 and their inverses.
- Magnetic refrigerators using paramagnetic salts attain  $\sim 0.2\text{K}$ . Adiabatic demagnetization attains mK temperatures.